

A point charge moving with constant speed.

① The fields in terms of the potentials,

$$\vec{B} = \vec{\nabla} \times \vec{A},$$

$$\vec{E} = -\vec{\nabla} \phi - \frac{\partial \vec{A}}{\partial t},$$

in the Lorenz gauge satisfy the non-homogeneous wave equations

$$\left(-\nabla^2 + \frac{1}{c^2} \frac{\partial^2}{\partial t^2}\right) \phi(\vec{r}, t) = \frac{1}{\epsilon_0} \rho(\vec{r}, t)$$

$$\left(-\nabla^2 + \frac{1}{c^2} \frac{\partial^2}{\partial t^2}\right) \vec{A}(\vec{r}, t) = \mu_0 \vec{J}(\vec{r}, t)$$

② Retarded Green's function:

$$\left(-\nabla^2 + \frac{1}{c^2} \frac{\partial^2}{\partial t^2}\right) \frac{1}{4\pi} G(\vec{r}-\vec{r}', t-t') = \delta^{(3)}(\vec{r}-\vec{r}') \delta(t-t')$$

with solution (satisfying causality)

$$G(\vec{r}-\vec{r}', t-t') = \frac{1}{|\vec{r}-\vec{r}'|} \delta\left(t-t' - \frac{1}{c} |\vec{r}-\vec{r}'|\right)$$

③ In terms of the Green's function we can solve

$$\phi(\vec{r}, t) = \frac{1}{4\pi \epsilon_0} \int d^3r' \int_{-\infty}^{+\infty} dt' G(\vec{r}-\vec{r}', t-t') \rho(\vec{r}', t')$$

$$\vec{A}(\vec{r}, t) = \frac{\mu_0}{4\pi} \int d^3r' \int_{-\infty}^{+\infty} dt' G(\vec{r}-\vec{r}', t-t') \vec{J}(\vec{r}', t')$$

④ Lienard-Wiechert potentials:

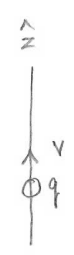
If we integrate over t' we obtain

$$\phi(\vec{r}, t) = \frac{1}{4\pi\epsilon_0} \int d^3x' \frac{1}{|\vec{r} - \vec{r}'|} \rho(\vec{r}', t - \frac{1}{c}|\vec{r} - \vec{r}'|)$$

$$\vec{A}(\vec{r}, t) = \frac{\mu_0}{4\pi} \int d^3x' \frac{1}{|\vec{r} - \vec{r}'|} \vec{J}(\vec{r}', t - \frac{1}{c}|\vec{r} - \vec{r}'|)$$

⑤ A point charge moving with constant speed is

described by



$$\vec{r}_a(t) = x_a(t) \hat{i} + y_a(t) \hat{j} + z_a(t) \hat{k} \\ = 0 \hat{i} + 0 \hat{j} + vt \hat{k}$$

$$\vec{v}_a(t) = \frac{d}{dt} \vec{r}_a(t) \\ = 0 \hat{i} + 0 \hat{j} + v \hat{k}$$

$$\rho(\vec{r}', t') = q \delta(x' - x_a(t')) \delta(y' - y_a(t')) \delta(z' - z_a(t')) \\ = q \delta(x') \delta(y') \delta(z' - vt')$$

$$\vec{J}(\vec{r}', t') = \vec{v}_a(t') \rho(\vec{r}', t') \\ = \hat{z} qv \delta(x') \delta(y') \delta(z' - vt')$$

⑥ Using ⑤ in ③ we have the scalar potential,

$$\begin{aligned} \phi(\vec{r}, t) &= \frac{1}{4\pi\epsilon_0} \int d^3\vec{r}' \int_{-\infty}^{+\infty} dt' G(\vec{r}-\vec{r}', t-t') \rho(\vec{r}', t') \\ &= \frac{q}{4\pi\epsilon_0} \int d^3\vec{r}' \int_{-\infty}^{+\infty} dt' \frac{\delta\left(t-t' - \frac{1}{c}|\vec{r}-\vec{r}'|\right)}{|\vec{r}-\vec{r}'|} \delta(x') \delta(y') \delta(z'-vt') \\ &= \frac{q}{4\pi\epsilon_0} \int_{-\infty}^{+\infty} dt' \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} d^3\vec{r}' \frac{\delta\left(t-t' - \frac{1}{c}\sqrt{(x-x')^2 + (y-y')^2 + (z-z')^2}\right)}{\sqrt{(x-x')^2 + (y-y')^2 + (z-z')^2}} \delta(x') \delta(y') \delta(z'-vt') \\ &= \frac{q}{4\pi\epsilon_0} \int_{-\infty}^{+\infty} dt' \frac{\delta\left(t-t' - \frac{1}{c}\sqrt{x^2 + y^2 + (z-vt')^2}\right)}{\sqrt{x^2 + y^2 + (z-vt')^2}} \end{aligned}$$

⑦ This involves an integral of the kind

$$\int_{-\infty}^{+\infty} dx \delta(F(x)) .$$

we shall digress to evaluate this integral.

⑧ Let us begin by evaluating

$$\int_{-\infty}^{+\infty} dx \delta(kx) = \frac{1}{k} \int_{-\infty}^{+\infty} dx' \delta(x') \quad (x' = kx, k > 0)$$

$$= \frac{1}{k}$$

$$\int_{-\infty}^{+\infty} dx \delta(kx) = \frac{1}{k} \int_{\infty}^{-\infty} dx' \delta(x') \quad (x' = kx, k < 0)$$

$$= -\frac{1}{k}$$

Together, we can write

$$\int_{-\infty}^{+\infty} dx \delta(kx) = \frac{1}{|k|},$$

which is stated, before completing the x' integral as.

$$\int_{-\infty}^{+\infty} dx \delta(kx) = \frac{1}{|k|} \int_{-\infty}^{+\infty} dx' \delta(x'),$$

or

$$\delta(kx) = \frac{\delta(x)}{|k|}.$$

⑨ Let us now evaluate

$$\begin{aligned} \int_{-\infty}^{+\infty} dx \delta((x-a_1)(x-a_2)) &= \int_{a_1-\delta}^{a_1+\delta} dx \delta((x-a_1)(x-a_2)) + \int_{a_2-\delta}^{a_2+\delta} dx \delta((x-a_1)(x-a_2)) \\ &= \int_{a_1-\delta}^{a_1+\delta} dx \delta((x-a_1)(a_1-a_2)) + \int_{a_2-\delta}^{a_2+\delta} dx \delta((a_2-a_1)(x-a_2)) \\ &= \frac{1}{|a_1-a_2|} + \frac{1}{|a_2-a_1|} \quad (\text{using } \textcircled{8}) \end{aligned}$$

Thus, we conclude

$$\delta((x-a_1)(x-a_2)) = \frac{\delta(x-a_1)}{|a_1-a_2|} + \frac{\delta(x-a_2)}{|a_2-a_1|}$$

⑩ We can immediately generalize the discussion in ⑨ to write

$$\int_{-\infty}^{+\infty} dx \delta((x-a_1) \dots (x-a_n)) = \sum_{i=1}^n \frac{\delta(x-a_i)}{\left| \frac{(x-a_1) \dots (x-a_n)}{(x-a_i)} \right|_{x=a_i}}$$

⑪ Next, we make the observation that if

$$f(x) = a_0 (x-a_1)(x-a_2)(x-a_3)$$

$$\begin{aligned} \text{then} \\ \frac{df}{dx} &= a_0 (x-a_2)(x-a_3) + a_0 (x-a_1)(x-a_3) + a_0 (x-a_1)(x-a_2) \\ &= \sum_{i=1}^3 \frac{f(x)}{(x-a_i)} \end{aligned}$$

(12) In general we have.

$$\frac{df}{dx} = \sum_r \frac{f(x)}{(x-a_r)}$$

a_r - roots of $f(x) = 0$.

(13) Using (12) in (10) we thus conclude that

$$\int_{-\infty}^{+\infty} dx \delta(F(x)) = \sum_r \frac{\delta(x-a_r)}{\left| \frac{\partial F}{\partial x} \Big|_{x=a_r} \right|}$$

(14) Using (13) in (6) we have.

$$\phi(\vec{r}, t) = \frac{q}{4\pi\epsilon_0} \int_{-\infty}^{+\infty} dt' \frac{1}{\sqrt{x^2 + y^2 + (z-vt')^2}} \sum_r \frac{\delta(t' - t_r)}{\left| \frac{\partial F}{\partial t'} \Big|_{t'=t_r} \right|}$$

where

$$F(t') = t - t' - \frac{1}{c} \sqrt{x^2 + y^2 + (z-vt')^2}$$

and t_r are the roots of the equation

$$F(t') = 0,$$

which is a quadratic equation - thus has two solutions.

(15) Using (14) we have.

$$\frac{\partial F}{\partial t'} = -1 + \frac{v}{c} \frac{(z-vt')}{\sqrt{x^2+y^2+(z-vt')^2}}$$

$$\left| \frac{\partial F}{\partial t'} \right| = 1 - \frac{v}{c} \frac{(z-vt')}{\sqrt{x^2+y^2+(z-vt')^2}}$$

(∵ $\frac{v}{c} < 1$
and $\frac{x}{\sqrt{a^2+x^2}} < 1$)

(16) Using (15) in (14)

$$\phi(r, t) = \frac{q}{4\pi\epsilon_0} \sum_{r=1}^2 \frac{1}{\sqrt{x^2+y^2+(z-vt_r)^2} - \frac{v}{c}(z-vt_r)}$$

$$= \frac{q}{4\pi\epsilon_0} \sum_{r=1}^2 \frac{1}{c(t-t_r) - \frac{v}{c}(z-vt_r)}$$

(using $F(t) = 0$)

$$= \frac{q}{4\pi\epsilon_0} \sum_{r=1}^2 \frac{1}{\left(t - \frac{v}{c} \frac{z}{c}\right) - t_r \left(1 - \frac{v^2}{c^2}\right)}$$

(17) Thus the problem reduces to finding the roots of

$$F(t) = t - t' - \frac{1}{c} \sqrt{x^2+y^2+(z-vt')^2} = 0$$

$$t_r^2 \left(1 - \frac{v^2}{c^2}\right) - 2 t_r \left(t - \frac{v}{c} \frac{z}{c}\right) + t^2 - \frac{x^2+y^2}{c^2} - \frac{z^2}{c^2} = 0$$

$$t_r = \frac{\left(t - \frac{v}{c} \frac{z}{c}\right) \pm \sqrt{\left(t - \frac{v}{c} \frac{z}{c}\right)^2 - \left(1 - \frac{v^2}{c^2}\right) \left(t^2 - \frac{x^2+y^2}{c^2} - \frac{z^2}{c^2}\right)}}{\left(1 - \frac{v^2}{c^2}\right)}$$

$$\begin{aligned}
 (18) \quad (\sqrt{\quad})^2 &= \sqrt{t^2 + \frac{v^2}{c^2} \frac{z^2}{c^2}} - 2 \frac{v}{c} \frac{z}{c} t - \left(1 - \frac{v^2}{c^2}\right) \left(t^2 - \frac{x^2 + y^2}{c^2} - \frac{z^2}{c^2}\right) \\
 &= -2 \frac{v}{c} \frac{z}{c} t + \frac{z^2}{c^2} + \frac{v^2}{c^2} t^2 + \left(1 - \frac{v^2}{c^2}\right) \frac{x^2 + y^2}{c^2} \\
 &= \frac{1}{c^2} (z - vt)^2 + \frac{1}{c^2} \left(1 - \frac{v^2}{c^2}\right) (x^2 + y^2)
 \end{aligned}$$

$z - vt \rightarrow$ retardation
 $1 - \frac{v^2}{c^2} \rightarrow$ Lorentz contraction.

(19) Using (18) in (17)

$$t_r = \frac{\left(t - \frac{v}{c} \frac{z}{c}\right) \pm \frac{1}{c} \sqrt{(x^2 + y^2) \left(1 - \frac{v^2}{c^2}\right) + (z - vt)^2}}{\left(1 - \frac{v^2}{c^2}\right)}$$

(20) To investigate which root contributes set $v = 0$.

$$t_r = t \pm \frac{1}{c} \sqrt{x^2 + y^2 + z^2}$$

Thus, retardation (or causal solution) corresponds to negative.

(21) Using (19) and (20) we conclude.

$$t_r = \frac{\left(t - \frac{v}{c} \frac{z}{c}\right) - \frac{1}{c} \sqrt{(x^2 + y^2) \left(1 - \frac{v^2}{c^2}\right) + (z - vt)^2}}{\left(1 - \frac{v^2}{c^2}\right)}$$

which can be rewritten as.

$$\left(t - \frac{v}{c} \frac{z}{c}\right) - t_r \left(1 - \frac{v^2}{c^2}\right) = \frac{1}{c} \sqrt{(x^2 + y^2) \left(1 - \frac{v^2}{c^2}\right) + (z - vt)^2}$$

(22) Using (21) in (16) we have, collecting only negative root,

$$\phi(\vec{r}, t) = \frac{1}{4\pi\epsilon_0} \frac{q}{\sqrt{(x^2+y^2)(1-\frac{v^2}{c^2}) + (z-vt)^2}}$$

→ $(z-vt)$ corresponds to retardation

→ $1-\frac{v^2}{c^2}$ factors are relativistic contributions — length contraction.

(23) Since the vector potential satisfies almost identical form we have.

$$\vec{A}(\vec{r}, t) = \frac{\mu_0}{4\pi} \frac{qv\hat{z}}{\sqrt{(x^2+y^2)(1-\frac{v^2}{c^2}) + (z-vt)^2}}$$

(24) We can now evaluate the fields.

$$\vec{E}(\vec{r}, t) = -\vec{\nabla}\phi(\vec{r}, t) - \frac{\partial}{\partial t}\vec{A}(\vec{r}, t)$$

(25) Using (23)

$$\begin{aligned} \frac{\partial}{\partial t}\vec{A}(\vec{r}, t) &= \frac{\mu_0}{4\pi} \hat{z} v^2 \frac{q(z-vt)}{[(x^2+y^2)(1-\beta^2) + (z-vt)^2]^{\frac{3}{2}}} \\ &= \frac{1}{4\pi\epsilon_0} \hat{z} \beta^2 \frac{q(z-vt)}{[(x^2+y^2)(1-\beta^2) + (z-vt)^2]^{\frac{3}{2}}} \end{aligned}$$

$$\beta = \frac{v}{c}$$

(26) Using (22)

$$\vec{\nabla} \phi(\vec{r}, t) = - \frac{q}{4\pi\epsilon_0} \frac{x(1-\beta^2)\hat{i} + y(1-\beta^2)\hat{j} + (z-vt)\hat{k}}{[(x^2+y^2)(1-\beta^2) + (z-vt)^2]^{\frac{3}{2}}}$$

(27) Using (25) and (26) in (24)

$$\vec{E}(\vec{r}, t) = \frac{q}{4\pi\epsilon_0} (1-\beta^2) \frac{x\hat{i} + y\hat{j} + (z-vt)\hat{k}}{[(x^2+y^2)(1-\beta^2) + (z-vt)^2]^{\frac{3}{2}}}$$

(28) $\vec{B}(\vec{r}, t) = \vec{\nabla} \times \vec{A}(\vec{r}, t)$

$$\begin{aligned} &= \vec{\nabla} \times \frac{\mu_0}{4\pi} \frac{qv\hat{z}}{\sqrt{(x^2+y^2)(1-\beta^2) + (z-vt)^2}} \\ &= - \frac{\mu_0}{4\pi} qv \hat{z} \times \vec{\nabla} \frac{1}{\sqrt{(x^2+y^2)(1-\beta^2) + (z-vt)^2}} \\ &= \frac{\mu_0}{4\pi} qv \hat{k} \times \frac{x(1-\beta^2)\hat{i} + y(1-\beta^2)\hat{j} + (z-vt)\hat{k}}{[(x^2+y^2)(1-\beta^2) + (z-vt)^2]^{\frac{3}{2}}} \\ &= \frac{\mu_0}{4\pi} qv (1-\beta^2) \frac{\sqrt{x^2+y^2} \hat{\phi}}{[(x^2+y^2)(1-\beta^2) + (z-vt)^2]^{\frac{3}{2}}} \end{aligned}$$

$$\hat{\phi} = \frac{x\hat{j} - y\hat{i}}{\sqrt{x^2+y^2}}$$

(29) Summary:

$$\phi(\vec{r}, t) = \frac{q}{4\pi\epsilon_0} \frac{1}{\sqrt{(x^2+y^2)(1-\beta^2) + (z-vt)^2}}$$

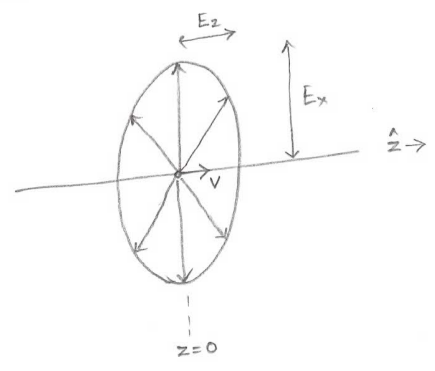
$$\vec{A}(\vec{r}, t) = \frac{1}{c} \frac{q}{4\pi\epsilon_0} \frac{v}{\sqrt{(x^2+y^2)(1-\beta^2) + (z-vt)^2}}$$

$$\vec{E}(\vec{r}, t) = (1-\beta^2) \frac{q}{4\pi\epsilon_0} \frac{x\hat{i} + y\hat{j} + (z-vt)\hat{k}}{[(x^2+y^2)(1-\beta^2) + (z-vt)^2]^{\frac{3}{2}}}$$

$$\vec{B}(\vec{r}, t) = \frac{1}{c} \beta (1-\beta^2) \frac{q}{4\pi\epsilon_0} \frac{(x\hat{j} - y\hat{i})}{[(x^2+y^2)(1-\beta^2) + (z-vt)^2]^{\frac{3}{2}}}$$

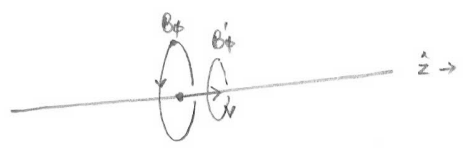
$$\hat{\phi} = \frac{x\hat{j} - y\hat{i}}{\sqrt{x^2+y^2}}$$

(30) Snapshot at $t=0$:



$$E_x = |\vec{E}(x, 0, 0, 0)| = \frac{1}{\sqrt{1-\beta^2}} \frac{q}{4\pi\epsilon_0} \frac{1}{x^2}$$

$$E_z = |\vec{E}(0, 0, z, 0)| = (1-\beta^2) \frac{q}{4\pi\epsilon_0} \frac{1}{z^2}$$



$$c B_\phi = |c \vec{B}(x, 0, 0, 0)| = \frac{\beta}{\sqrt{1-\beta^2}} \frac{q}{4\pi\epsilon_0} \frac{1}{x^2}$$

$$c B'_\phi = |c \vec{B}(x, 0, z, 0)| = \beta (1-\beta^2) \frac{q}{4\pi\epsilon_0} \frac{x}{[x^2(1-\beta^2) + z^2]^{\frac{3}{2}}}$$