

Modified Bessel functions

① For a planar geometry we have

$$G(\vec{r}, \vec{r}') = \frac{1}{\epsilon_0} \int \frac{d^2 k_{\perp}}{(2\pi)^2} e^{i\vec{k}_{\perp} \cdot (\vec{r} - \vec{r}')_{\perp}} \frac{1}{2k_{\perp}} e^{-k_{\perp} |z - z'|}$$

$$= \frac{1}{4\pi \epsilon_0} \int_0^{\infty} dk_{\perp} J_0(k_{\perp} P) e^{-k_{\perp} |z - z'|}$$

$\vec{P} = (\vec{r} - \vec{r}')_{\perp}$

② Observe the integral

$$\int_{-\infty}^{+\infty} \frac{dk_z}{2\pi} \frac{e^{ik_z(z-z')}}{k_z^2 + k_{\perp}^2} = \frac{1}{2k_{\perp}} e^{-k_{\perp} |z-z'|}$$

→ we contour integrate to evaluate this HW.

③ Using ② in ①

$$G(\vec{r}, \vec{r}') = \frac{1}{2\pi \epsilon_0} \int_0^{\infty} k_{\perp} dk_{\perp} J_0(k_{\perp} P) \int_{-\infty}^{+\infty} \frac{dk_z}{2\pi} \frac{e^{ik_z(z-z')}}{k_z^2 + k_{\perp}^2}$$

$$= \frac{1}{2\pi \epsilon_0} \int_{-\infty}^{+\infty} \frac{dk_z}{2\pi} e^{ik_z(z-z')} \boxed{\int_0^{\infty} k_{\perp} dk_{\perp} \frac{J_0(k_{\perp} P)}{k_{\perp}^2 + k_z^2}}$$

$K_0(k_z P)$

④ Modified Bessel function is defined as

$$K_0(t) = \int_0^{\infty} s ds \frac{J_0(s)}{s^2 + t^2} \quad 0 < t < \infty$$

⑤ Thus, we have.

$$G(\vec{r}, \vec{r}') = \frac{1}{4\pi\epsilon_0} \int_0^\infty dk_\perp J_0(k_\perp \rho) e^{-k_\perp |z-z'|} \rightarrow \text{planar geometry}$$

$$= \frac{1}{2\pi\epsilon_0} \int_{-\infty}^\infty \frac{dk_z}{2\pi} e^{ik_z(z-z')} K_0(k_z \rho) \rightarrow \text{cylindrical geometry}$$

$$\begin{aligned} \text{⑥ } \frac{1}{2\pi\epsilon_0} \int_{-\infty}^0 \frac{dk_z}{2\pi} e^{ik_z(z-z')} K_0(k_z \rho) &= \frac{1}{2\pi\epsilon_0} \int_0^\infty \frac{dk'_z}{2\pi} e^{-ik'_z(z-z')} K_0(-k'_z \rho) \quad k_z = -k'_z \\ &= \frac{1}{2\pi\epsilon_0} \int_0^\infty \frac{dk_z}{2\pi} e^{-ik_z(z-z')} K_0(k_z \rho) \end{aligned}$$

Thus we have.

$$G(\vec{r}, \vec{r}') = \frac{1}{2\pi\epsilon_0} \int_0^\infty \frac{dk_z}{2\pi} \left[e^{ik_z(z-z')} + e^{-ik_z(z-z')} \right] K_0(k_z \rho)$$

$$= \frac{1}{4\pi\epsilon_0} \frac{2}{\pi} \int_0^\infty dk_z \cos[k_z(z-z')] K_0(k_z \rho)$$

which brings out the oscillatory nature in the z-direction more explicitly.

$$\begin{aligned}
(7) \quad K_0(k_z P) &= \int_0^\infty k_\perp dk_\perp \frac{J_0(k_\perp P)}{k_\perp^2 + k_z^2} \\
&= \int_0^\infty k_\perp dk_\perp \int_0^{2\pi} \frac{d\alpha}{2\pi} e^{i k_\perp P \cos \alpha} \frac{1}{k_\perp^2 + k_z^2} \\
&= 2\pi \int \frac{d^2 k_\perp}{(2\pi)^2} e^{i \vec{k}_\perp \cdot (\vec{r} - \vec{r}')_\perp} \frac{1}{k_\perp^2 + k_z^2}
\end{aligned}$$

(8) The above form immediately suggests

$$\begin{aligned}
(-\nabla_\perp^2 + k_z^2) \frac{1}{2\pi} K_0(k_z P) &= \int \frac{d^2 k_\perp}{(2\pi)^2} (-\nabla_\perp^2 + k_z^2) e^{i \vec{k}_\perp \cdot (\vec{r} - \vec{r}')_\perp} \frac{1}{k_\perp^2 + k_z^2} \\
&= \int \frac{d^2 k_\perp}{(2\pi)^2} (-(i k_\perp)^2 + k_z^2) e^{i \vec{k}_\perp \cdot (\vec{r} - \vec{r}')_\perp} \frac{1}{k_\perp^2 + k_z^2} \\
&= \int \frac{d^2 k_\perp}{(2\pi)^2} e^{i \vec{k}_\perp \cdot (\vec{r} - \vec{r}')_\perp} \\
&= \delta^{(2)}(\vec{r}_\perp - \vec{r}'_\perp)
\end{aligned}$$

Thus, we have the Green's function equation:

$$[-\nabla_\perp^2 + k_z^2] \frac{1}{2\pi} K_0(k_z P) = \delta^{(2)}(\vec{r}_\perp - \vec{r}'_\perp)$$

⑨

Using

$$\delta^{(2)}(\vec{r}_1 - \vec{r}'_1) = \frac{\delta(r - r')}{r} \delta(\phi - \phi')$$

$$\vec{r}_1 = (r \cos \phi, r \sin \phi)$$

$$\vec{r}'_1 = (r' \cos \phi', r' \sin \phi')$$

$$\vec{P} = (r \cos \phi - r' \cos \phi', r \sin \phi - r' \sin \phi')$$

$$P = \sqrt{r^2 + r'^2 - 2rr' \cos(\phi - \phi')}$$

$$\tan \alpha = \frac{r \sin \phi - r' \sin \phi'}{r \cos \phi - r' \cos \phi'}$$

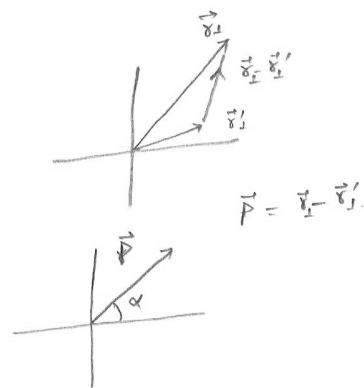
In the limit $\vec{r}_1 \rightarrow \vec{r}'_1$, $r \rightarrow r'$

$$P \rightarrow r \sqrt{2 - 2 \cos(\phi - \phi')} = 2r \sin\left(\frac{\phi - \phi'}{2}\right) \rightarrow 0$$

$$\tan \alpha \rightarrow -\frac{r(\phi - \phi')}{r \frac{\phi^2}{2} - \frac{\phi'^2}{2}} = -\frac{2}{\phi + \phi'} = -\frac{1}{\phi} \quad ?$$

Thus, argue that, which is intuitively clear

$$\delta^{(2)}(\vec{r}_1 - \vec{r}'_1) = \frac{\delta(r - r')}{r} \delta(\phi - \phi') = \frac{\delta(P)}{P} \frac{1}{2\pi}$$



⑩ Using ⑨ in ⑧ we have.

$$[-\nabla_x^2 + k_z^2] \frac{1}{2\pi} K_0(k_z P) = \frac{\delta(P)}{P} \frac{1}{2\pi}$$

$$\left[-\frac{1}{P} \frac{d}{dP} P \frac{d}{dP} + k_z^2\right] K_0(k_z P) = \frac{\delta(P)}{P}$$

$$\left[-\frac{1}{t} \frac{d}{dt} t \frac{d}{dt} + 1\right] K_0(t) = \frac{\delta(t)}{t}$$

$$\left[-\frac{1}{t} \frac{d}{dt} t \frac{d}{dt} + 1\right] K_0(t) = 0 \quad t \neq 0.$$

⑪ Replacing $t \rightarrow it$

$$\left[\frac{1}{t} \frac{d}{dt} t \frac{d}{dt} + 1\right] K_0(it) = 0$$

which reveals that $J_0(t) = K_0(it)$.

⑫

$$\left[-\frac{1}{P} \frac{d}{dP} P \frac{d}{dP} + k_z^2\right] K_0(k_z P) = \left[-\frac{1}{P} \frac{d}{dP} P \frac{d}{dP} + k_z^2\right] \int_0^\infty k_\perp dk_\perp \frac{J_0(k_\perp P)}{k_\perp^2 + k_z^2}$$

$$= \int_0^\infty k_\perp dk_\perp \left[-\frac{1}{P} \frac{d}{dP} P \frac{d}{dP} + k_z^2\right] J_0(k_\perp P) \frac{1}{k_\perp^2 + k_z^2}$$

$$= \int_0^\infty k_\perp dk_\perp J_0(k_\perp P)$$

$$= \frac{\delta(P)}{P}$$

$$\int_0^\infty k_\perp dk_\perp J_0(k_\perp P) J_0(k_\perp P') = \frac{\delta(P P')}{P}$$

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$$\left[\overbrace{-\nabla_{\perp}^2} + k_z^2 \right] K_0(k_z P) = \frac{\delta(P)}{P}$$

Using divergence theorem in two dimension.

$$-\int d\vec{s} \cdot \vec{\nabla}_{\perp} K_0(k_z P) = 2\pi$$

↓
circle.

$$-2\pi P \frac{d}{dP} K_0(k_z P) = 2\pi$$

$$\frac{d}{dt} K_0(t) = -\frac{1}{t}$$

$$K_0(t) = -\ln t + \text{const}$$

$$= \ln \frac{2}{t} - \gamma$$

Euler's constant
 $\gamma = 0.577$